$(2.6 \pm 0.7) \times 10^8 \text{ M}^{-1} \text{ s}^{-1}$ within experimental uncertainty of the k_{co} reported for eq 3.

These results indicate that $H_2IrCl(CO)(PPh_3)_2$ first undergoes photodissociation of CO (eq 5) followed by elimination of H_2 from

$$H_2IrCl(CO)(PPh_3)_2 \xrightarrow{h\nu} H_2IrCl(PPh_3)_2 + CO$$
 (5)

the resulting pentacoordinated intermediate (eq 6) to give the

$$H_2IrCl(PPh_3)_2 \xrightarrow{k_5}_{k_4} IrCl(PPh_3)_2 + H_2$$
(6)

"IrCl(PPh₃)₂" transient formed directly via flash photolysis of trans-IrCl(CO)(PPh₃)₂. This view⁸ contrasts to the proposal⁹ that the dihydride photoelimination occurs by a single concerted step but is consistent with theoretical arguments⁴ and several experimental observations^{10,11} that reductive elimination from d^6 complexes often occurs much more readily after ligand dissociation from the original hexacoordinate species to give a pentacoordinate intermediate. Given that formation of IrCl(PPh₃)₂ was complete within the lifetime of the flash (20 μ s), a lower limit for k_6 can be estimated as 5×10^4 s⁻¹. Thus, we conclude that dissociation of CO accelerates dihydrogen elimination by at least 9 orders of magnitude. Notably, this rate acceleration occurs despite the dissociation of the π -acid CO which would be expected to favor the lower oxidation state of the metal center. Another notable feature is that the equilibrium indicated by eq 6 apparently lies largely to the right in contrast to the observation that under similar conditions the equilibrium for the Rh(I) analogue must lie to the left.7

Preliminary investigations of the Ir(III) species HIrCl- $(CO)(PPh_2C_6H_4)(PPh_3)$,¹² the orthometalated isomer of Vaska's compound, as well as of the dinitrogen Ir(I) complex trans- $IrCl(N_2)(PPh_3)_2^{13}$ demonstrate that flash photolysis leads in both cases to immediate appearance of a transient spectrum qualitatively the same as that attributed to " $IrCl(PPh_3)_2$ ". In the former case the transient decays by a second-order pathway (P_{CO} dependent, i.e., eq 3) to form trans-IrCl(CO)(PPh₃)₂ as the photoproduct. Again it appears that the starting complex has undergone CO photodissociation to give a pentacoordinate intermediate followed by H/aryl elimination to form IrCl(PPh₃), within the 20-µs lifetime of the flash. Thus, reductive elimination from this pentacoordinate intermediate must be at least 8 orders of magnitude faster than the rate of about $3 \times 10^{-4} \text{ s}^{-1}$ (70 °C) we have measured for the thermal reaction of HIrCl(CO)- $(PPh_2C_6H_4)(PPh_3)$ to give trans-IrCl(CO)(PPh_3)₂.¹⁴ For $IrCl(N_2)(PPh_3)_2$, photolabilization of N₂ to give $IrCl(PPh_3)_2$ is irreversible and, in the absence of the other reactants, this reactive intermediate apparently undergoes internal orthometalation of a triphenylphosphine. This is accompanied by slow absorbance decreases in the 340-550-nm range with isosbestic points at 460 and 334 nm consistent with formation of Ir(III) products. The

proton NMR spectrum of the product solutions from the continuous photolysis of $IrCl(N_2)(PPh_3)_2$ in C_6D_6 under otherwise analogous conditions displays a -22.5 ppm resonance indicating formation of an iridium hydride.¹⁵ These investigations will be described in detail in subsequent publications.

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Diagonal Peak Suppression in 2D-NOE Spectra

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The 2D correlation NMR experiments COSY¹ and NOESY² are standard techniques for study of molecular structure in solution. COSY, which is based on coherence transfer via the J coupling, has been the progenitor of a myriad of other pulse sequences for specific applications.³⁻⁵ Particularly useful have been filtered COSY pulse sequences,⁶ which are conceptually derived from multiple-quantum selective-excitation experiments.^{7,1} Double-quantum filtered COSY especially is frequently employed to suppress single-quantum peaks, improving the dynamic range and the resolution close to the diagonal.9 NOESY, which detects through-space magnetization-transfer via direct dipolar relaxation, frequently suffers from a more severe dynamic range problem than COSY. While use has recently been made of coherence transfer via the J coupling to selectively observe certain resonances in NOESY spectra, 10 until now there has not been a general method for removing the diagonal. Thus, cross-peaks between nuclei with similar chemical shifts, which fall close to the diagonal, are obscured. Furthermore, the spectra are also often marred by strong solvent lines and t_1 noise. In this paper we describe a magneti-

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⁽⁸⁾ The following observation argues against a sequential two-photon process involving initial H₂ photolabilization to generate $IrCl(PPh_3)_2$ followed by secondary photolysis of this product to give " $IrCl(PPh_3)_2$ ". The relative pulse intensity required to generate the same concentration of the latter transient was 5 times larger when the initial substrate was *trans*-IrCl(CO)(PPh_3)₂ (under argon) than when H₂IrCl(CO)(PPh_3)₂ (under H₂) was the initial substrate. An alternative mechanism should be considered, namely, the possibility that both CO and H₂ are dissociated from the excited state of H₂IrCl(CO)(PPh_3)₂. At present such a mechanism, although unprecedented, cannot be differentiated from the stepwise pathway proposed in eq 5 and 6.

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Figure 1. (a) ¹H NOESY spectrum (500 MHz, 35 °C) of a 5% solution of *N*,*N*-dimethylacetamide in acetonitrile-*d*₃. The diagonal peaks due to the amide syn- and anti-CH₃ groups, at (ω_A, ω_A) and (ω_B, ω_B) are associated with cross-peaks at (ω_A, ω_B) and (ω_B, ω_A) arising from intramolecular chemical exchange of these residues. The cross-peak intensity amounts to approximately 30% of that of the diagonal. The signal at (ω_C, ω_C) derives from the (nonexchanging) COCH₃ group. The spectrum was acquired with sweep widths of 1 kHz in both dimensions and a t_m of 500 ms which included a 20-ms homospoil pulse. (b) ¹H spectrum of the same sample obtained with the diagonal-suppression technique. Note that the diagonal lines now appear as a horizontal row at $\omega_1 = 0$, with an intensity of about 3% of the diagonal lines in the NOESY spectrum. This spectrum was obtained under identical conditions with those used for (a), except that the sweep width in ω_1 was 2 kHz.

zation-transfer experiment which appears to offer a general solution to the dynamic-range problem in NOESY spectra and in which diagonal peaks and t_1 noise are suppressed, considerably improving spectral quality.

Our experiment consists of four pulses, a mixing (t_m) period, *twin* evolution (t_1) periods, and an acquisition (t_2) period as shown below:

$$P1(90^{\circ},\phi_1)-t_1-P2(90^{\circ},\phi_2)-t_m-P3(90^{\circ},\phi_3)-t_1-P4(90^{\circ},\phi_4)-ACQ$$

The pulse sequence can easily be understood by first considering a thought experiment in which $\phi_4 = 180^\circ + \phi_1$, $\phi_2 = \phi_3$, and t_m is infinitesimal. During the first t_1 period the individual spins undergo forward precession. If nothing occurs during t_m , then P2 and P3 constitute a 180° refocusing pulse, so that precession during the second t_1 brings the magnetization vectors back perpendicular to ϕ_1 , resulting in no signal during acquisition. Suppose, however, that some exchange of magnetization occurs during the infinitesimal t_m . Then precession in the second t_1 will be at a different frequency from that in the first, and the magnetization will no longer echo at P4; rather, it will lie somewhere in the XY plane, and P4 will conserve the component along ϕ_4 . The signal from the cross-relaxing species will be proportional to sin ($\omega_{\rm B}$ – $(\omega_A)t_1$, where ω_B and ω_A are the precession frequencies in the second and first t_1 , respectively. Fourier transformed and phased, this gives an antisymmetric spectrum with a positive peak at $\omega_{\rm B} - \omega_{\rm A}$ and a negative peak at $\omega_A - \omega_B$, similar to a zero-quantum experiment.11-13

For $t_m > 0$, however, this simple scheme does not work because P2 rotates the component of the complex magnetization perpendicular to ϕ_2 to $\pm Z$, leaving the other component invariant and



Figure 2. (Top) ¹H 1D-NMR spectrum (500 MHz, 30 °C) of the DNA dodecamer AATTbR⁵CGbR⁵CGAATT as a 1.4 mM solution in D₂O. (left) Contour plot of the 3.7-5.2 ppm region of the NOESY spectrum of the dodecamer, obtained with $t_m = 320$ ms, including a 20-ms homospoil, a recycle delay of 2 s, and sweep widths in ω_1 and ω_2 of 4 kHz. Note the strong diagonal ridge, the intense t_1 noise at $\omega_2 = 4.68$ ppm, and the crowded 4.0-4.25 ppm region. (Right) Contour plot of the same region of the 500-MHz diagonal-suppressed 2D-NOE spectrum of the same sample. Pairs of cross-peaks are connected by diagonals of slope $\Delta\omega_1/\Delta\omega_2 = 2$: this is demonstrated in the figure for the terminal 4' to 5',5'' lines. Only positive contours are plotted. Note the lower level of t_1 noise and the absence of significant diagonal intensity. The improved resolution of our methods is best appreciated by comparing the terminal 5',5" proton signals (circled). Cross-relaxation peaks are resolved only in the diagonal-suppressed spectrum. Again the two 2D spectra were obtained under identical conditions except that the sweep width in ω_2 of the diagonal-suppressed spectrum was 8 kHz.

in the XY plane. The latter component dephases much faster than the former, since it is destroyed by T_2 processes as well as T_1 ; therefore, the magnetization after P3 is no longer a 180° rotation of that before P2. We can, however, preserve the real and imaginary components separately along Z, destroying the other with a homospoil pulse. This is done by conducting two identical experiments in which $(\phi_2)_2 = (\phi_2)_1 + 90^\circ$. If then we make $(\phi_3)_1$ = $(\phi_2)_1$, as before, and $(\phi_3)_2 = -(\phi_3)_2$, and add the resulting acquisitions, the summed magnetization after P3 will once again be a 180° rotation of that before P2. Thus, if no cross-relaxation occurs, an echo will again be formed at P4 and the non-crossrelaxing magnetization returned to Z; while the cross-relaxing peaks will again be modulated by sin $(\omega_{\rm B} - \omega_{\rm A})t_1$, exactly as in the thought experiment. As only the component of the magnetization parallel to ϕ_4 contains any information, the perpendicular component is discarded after the first transform, thus eliminating half the noise. This partially mitigates the loss in sensitivity over the NOESY experiment.

Figure 1 compares spectra obtained with our method and NOESY from a standard sample, N,N-dimethylacetamide, in which the amide methyl groups rapidly exchange with each other. In the NOESY spectrum (Figure 1a) this is revealed by crosspeaks between the two downfield proton resonances at (ω_A, ω_B) and (ω_B, ω_A) , with about 30% of the intensity of the diagonal, nonexchanged signals (ω_A, ω_A) , (ω_B, ω_B) , and (ω_C, ω_C) . In the spectrum obtained with the new technique the diagonal peaks appear at $\omega_1 = 0$ and are less than 3% of their intensity in the NOESY experiment. Methyl resonances which undergo exchange during t_m give positive peaks at $(\omega_B - \omega_A, \omega_A)$ and $(\omega_A - \omega_B, \omega_B)$, and negative ones at $(\omega_A - \omega_B, \omega_A)$ and $(\omega_B - \omega_A, \omega_B)$. The efficiency of diagonal suppression is quite apparent.

In Figure 2 we apply the pulse sequence to a more challenging system, a 1.4 mM D_2O solution of the DNA dodecamer AATTbR⁵CGbR⁵CGAATT, containing brominated cytosines, again comparing it to a NOESY spectrum taken under identical conditions. The contour plots depict the highly overlapped 3.7–5.2 ppm region, which contains the 48 3', 4', 5', and 5" proton res-

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onances. The improved resolution of peaks close to the diagonal is apparent, as is the suppression of t_1 noise associated with the HDO signal at 4.68 ppm. Particularly impressive are the cross-peaks between the terminal 5' and 5'' protons at $\omega_2 = 3.7$ ppm. These are resolved neither in the 1D spectrum nor via NOESY. It should be pointed out that all of the cross-peaks in the region are relatively intense; where weakly coupled protons of similar chemical shift exist, the advantage of the diagonalsupression pulse sequence over NOESY will be amplified. In a forthcoming paper¹⁴ we shall discuss in more depth the pulse sequence and its applications in both solution and the solid state.

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Synthesis and Deuteration of $(\eta$ -Thiophene)Ru $(\eta$ -C₅H₅)⁺: A Model for Adsorption and Deuterium Exchange of Thiophene on Hydrodesulfurization Catalysts

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While the hydrodesulfurization (HDS) of thiophene on heterogeneous catalysts has been studied extensively,² the complexity of this reaction has left most important mechanistic questions unanswered. For other heterogeneously catalyzed reactions, important mechanistic concepts have been developed by studying homogeneous reactions of model compounds.³ Few model studies of thiophene HDS have been reported.⁴ Indeed, it has not been established whether thiophene adsorbs to the surface by coordination to a metal site through the sulfur atom only⁵ or via the sulfur and the unsaturated carbon atoms (i.e., as a π -complex).⁶ Recently, Harris and Chianelli⁷ concluded, on the basis of SCF-X α molecular orbital calculations of metal sulfides, that the initial interaction of thiophene with second-row, catalytically active transition metals was most likely via the S atom only. They noted that the preference for S atom coordination by second-row metals was important for the higher HDS activity of these metals, in-



Figure 1.

cluding Ru. Supporting the conclusion that second-row metals preferred S atom coordination were the facts that there were no known π -thiophene complexes of the second- and third-row transition metals and that the only reasonably well-characterized S-bound thiophene complex was of Ru, namely, Ru(NH₃)₅- $(SC_4H_4)^{2+8.9}$ In this paper, we report the first example of a π -thiophene complex, $(\eta - C_4 H_4 S) Ru(\eta - C_5 H_5)^+(1)$, of a secondor third-row transition metal (Figure 1). The thiophene is strongly coordinated and readily exchanges its 2,5-protons in the presence of bases or Al₂O₃; this latter observation suggests an explanation for the known exchange of these protons when thiophene and D₂ are passed over HDS catalysts.

The $[(\eta - C_4H_4S)Ru(\eta - C_5H_5)]BF_4$ (1) complex was prepared by refluxing $(\eta - C_5 H_5) Ru(PPh_3)_2 Cl^{10}$ (1.00 g, 1.38 mmol), thiophene (20 mL), AgBF₄ (0.290 g, 1.52 mmol) in 10 mL of MeOH for 72 h under N_2 . Isolation of 1 was accomplished by removing the volatiles from the reaction mixture in vacuo, extracting the residue with CH₂Cl₂, and precipitating the product by slowly adding Et_2O . Successive recrystallizations yielded the pure air- and water-stable 1 as a pale brown powder in 60% yield.¹¹ The π -thiophene ligand is strongly bound to the Ru as shown by its slow rate of displacement (only 33% after 4.5 h) from 1 (10 mg) by (n-Bu)₃P (7 equiv) in acetone (0.35 mL) at room temperature.

The stability of 1 suggests that thiophene could adsorb to HDS catalysts in the π -bonded form. This form may account for the observed exchange of the 2,5-protons of thiophene when it is passed with D₂ over several catalysts (e.g., Mo/Al₂O₃, 12 Mo-Co/Al₂O₃, 12 MoS_2 ,¹³⁻¹⁵ and M_xMoS_2 ¹⁵). From ¹H NMR studies we find that the 2,5-protons of the thiophene ligand in 1 readily undergo exchange in the presence of bases. Thus, the 2,5-protons in 1 (0.003 mmol) completely exchange with deuterium (CD_3OD) solvent 0.35 mL) in the presence of KOH (0.01 mmol) in less than 4 min.¹⁶ Deuteration at the 2,5-positions was established by displacing the exchanged thiophene from 1 with tert-butyl isocyanide and noting the almost complete disappearance of the 2,5-proton resonance of free thiophene.¹⁷ There was no evidence for exchange of the 3,4-protons. Treatment of 1 (0.017 mmol) in CD₃OD (0.35 mL) with Et₃N (0.04 mmol) resulted in approximately 50% exchange after 20 min. No exchange is observed

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